Dynamics of Systems

CTB 2300

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CTB 2300 Dynamics of Systems Faculty of Civil Engineering and Geosciences 2022 Delft, The Netherlands Department: Hydraulic Engineering Research group: Offshore Engineering Room: 3.71 E-mail: <u>H.Hendrikse@tudelft.nl</u> Web: tinyurl.com/hhendrikse



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Contents of Lecture 12

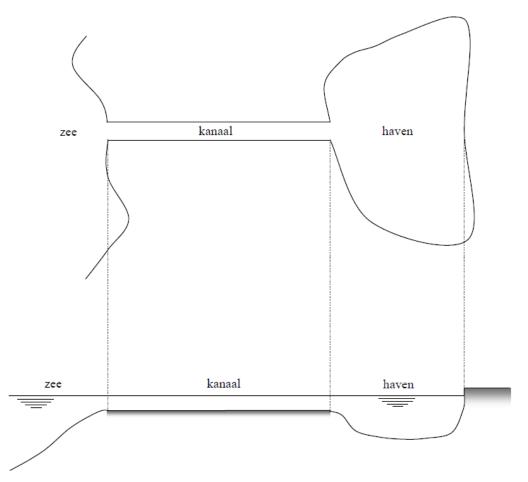
- 1. Equation of motion for the water level in a harbour connected with Sea by a canal
- 2. Analogy with the mechanical oscillations
- **3.** Analysis of the hydraulic systems based on the solutions obtained for the mechanical systems
- 4. Other examples of oscillatory dynamical systems

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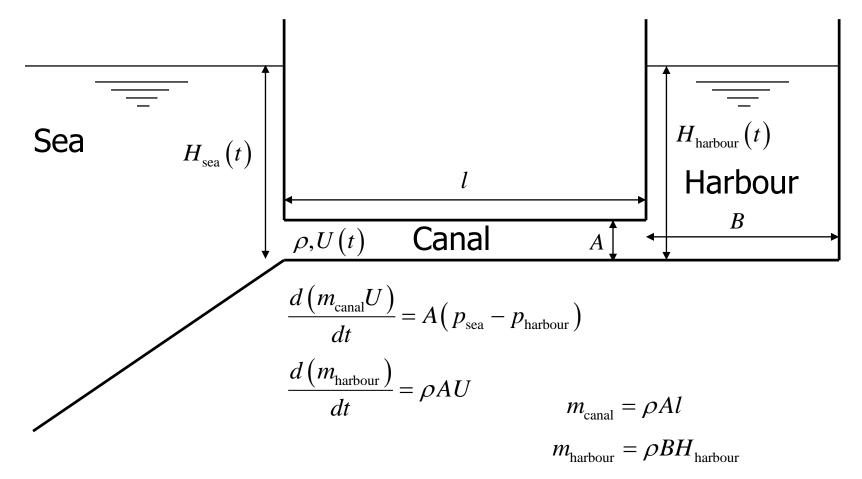


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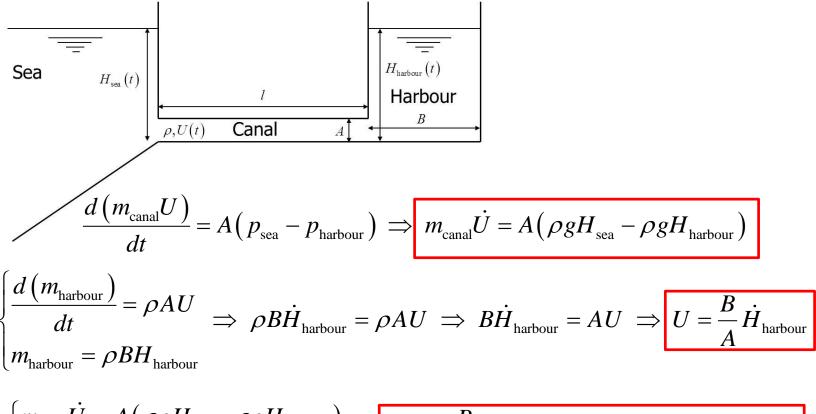


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$$m_{\text{canal}} U = A \left(\rho g H_{\text{sea}} - \rho g H_{\text{harbour}} \right) \implies \left(\rho A l \right) \frac{B}{A} \ddot{H}_{\text{harbour}} = A \left(\rho g H_{\text{sea}} - \rho g H_{\text{harbour}} \right)$$
$$m_{\text{canal}} = \rho A l$$

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$$(\rho Al)\frac{B}{A}\ddot{H}_{\text{harbour}} = A(\rho g H_{\text{sea}} - \rho g H_{\text{harbour}}) \implies \rho Bl\ddot{H}_{\text{harbour}} + \rho Ag H_{\text{harbour}} = \rho Ag H_{\text{sea}}$$

$$\rho B \ddot{H}_{\text{harbour}} + \rho A \frac{g}{l} H_{\text{harbour}} = \rho A \frac{g}{l} H_{\text{sea}}$$

 ρB - mass of water in the harbour per unit height

 $\rho A \frac{g}{l}$ - force from the harbour on the water in the canal per unit height (equivalent spring constant)

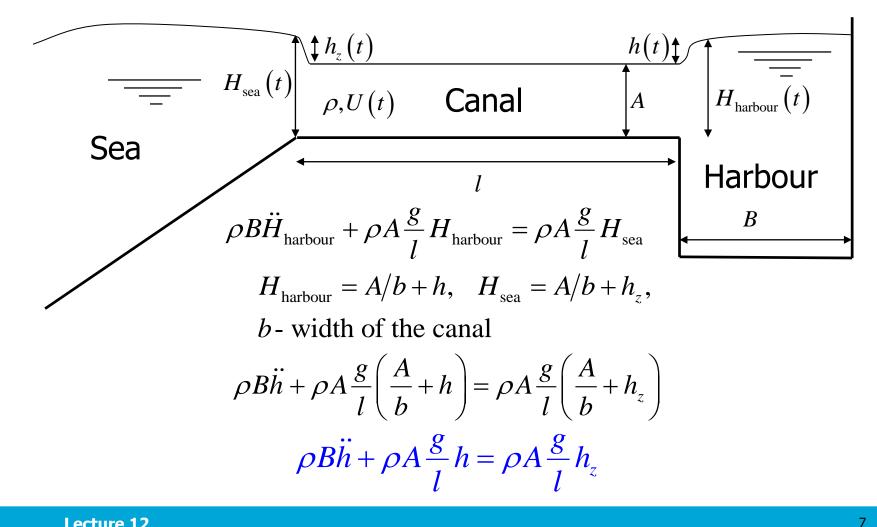
 $\rho A \frac{g}{l}$ - force from the sea on the water in the canal per unit height

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Analogy with the mass-spring system subject to a force

$$\rho B\ddot{h} + \rho A \frac{g}{l}h = \rho A \frac{g}{l}h_z \quad \text{Equations of motion} \qquad m\ddot{u} + ku = F_u$$

$$\rho B \leftrightarrow m \quad (\text{mass per unit height} \leftrightarrow \text{mass})$$

$$\rho A \frac{g}{l} \leftrightarrow k \quad (\text{equivalent stiffness} \leftrightarrow \text{stiffness})$$

$$\rho A \frac{g}{l}h_z \leftrightarrow F_u \qquad (\text{force} \leftrightarrow \text{force})$$

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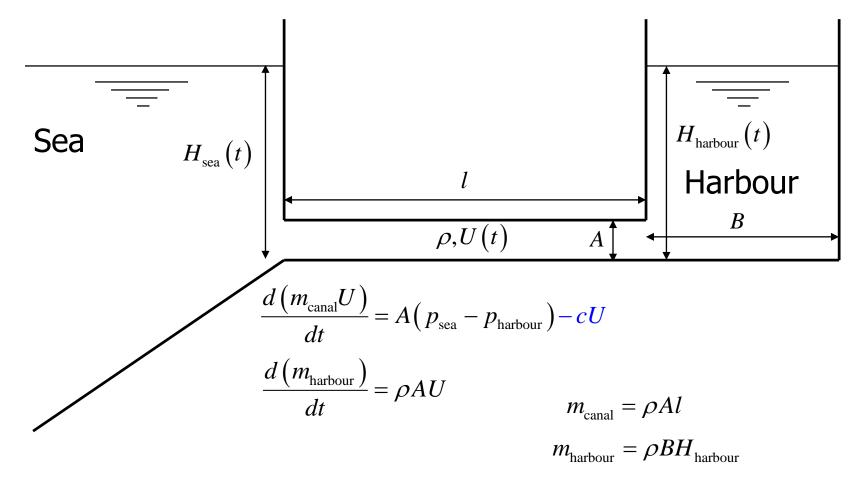
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 $\rho A \frac{s}{l} h_z \leftrightarrow F_u$



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Effect of the friction with the canal bottom



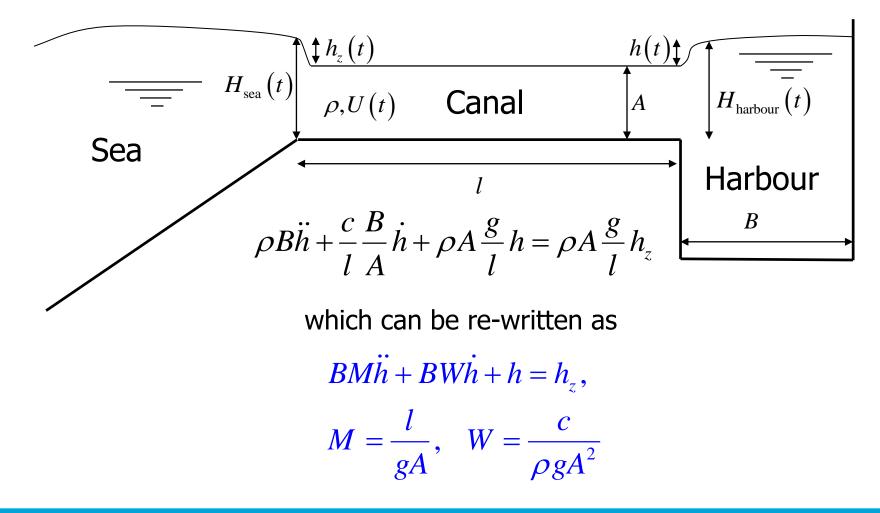
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Effect of the friction with the canal bottom



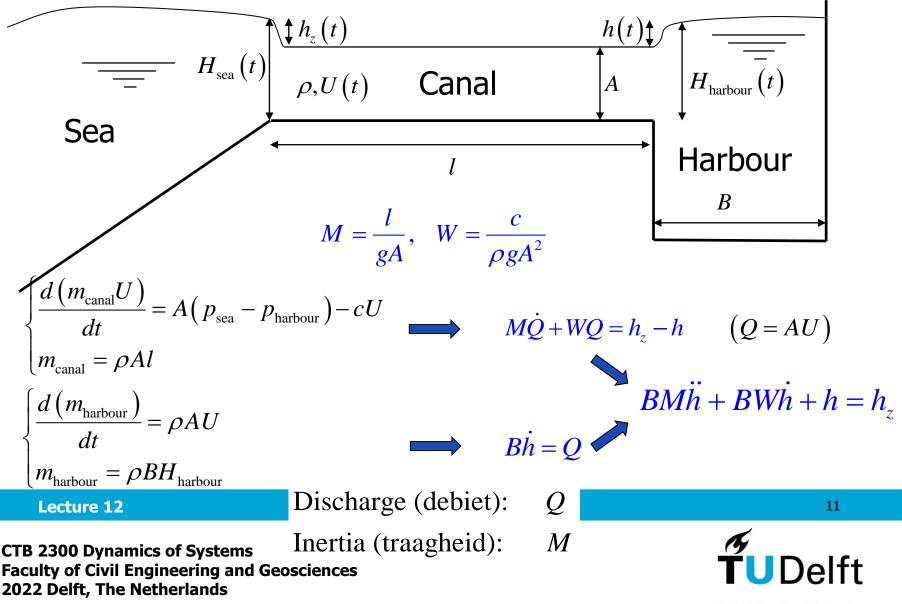
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Some specific hydraulic definitions



How to analyze?

$$BM\ddot{h} + BW\dot{h} + h = h_z$$

$$\downarrow$$

$$\ddot{h} + \frac{W}{M}\dot{h} + \frac{1}{BM}h = \frac{1}{BM}h_z$$

$$\downarrow$$

$$\ddot{h} + 2\zeta\omega_0\dot{h} + \omega_0^2h = f(t)$$

$$\omega_0^2 = \frac{1}{BM}, \ \zeta = \frac{W}{2}\sqrt{\frac{B}{M}}, \ f(t) = \frac{h_z(t)}{BM}$$

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We can simply use all formulas obtained for the massspring-dashpot system

$$\ddot{h} + 2\zeta \omega_0 \dot{h} + \omega_0^2 h = f(t)$$
$$\omega_0^2 = \frac{1}{BM}$$
$$\zeta = \frac{W}{2} \sqrt{\frac{B}{M}}$$
$$f(t) = \frac{h_z(t)}{BM}$$
hydraulic system

$$\ddot{u} + 2\zeta \omega_0 \dot{u} + \omega_0^2 u = f(t)$$
$$\omega_0^2 = \frac{k}{m}$$
$$\zeta = \frac{c}{2\sqrt{km}}$$
$$f(t) = \frac{F(t)}{m}$$
mechanical system

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Example 1: free vibration of the hydraulic system

As derived in Lecture 6, slide 19, the sub-critically damped free vibration of a mechanical system is described by

$$u(t) = \exp(-\zeta\omega_0 t) \left(u_0 \cos(\omega_1 t) + \frac{v_0 + \zeta\omega_0 u_0}{\omega_1} \sin(\omega_1 t) \right)$$

Therefore, using the expressions given on the previous slide, we find the following expression for the hydraulic system

$$h(t) = \exp\left(-\zeta\omega_0 t\right) \left(h(0)\cos\left(\omega_1 t\right) + \frac{\dot{h}(0) + \zeta\omega_0 h(0)}{\omega_1}\sin\left(\omega_1 t\right)\right)$$
$$\omega_0^2 = \frac{1}{BM}, \ \zeta = \frac{W}{2}\sqrt{\frac{B}{M}}, \ \omega_1 = \omega_0\sqrt{1-\zeta^2} = \frac{1}{\sqrt{BM}}\sqrt{1-\frac{W^2}{4}\frac{B}{M}}$$

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Example 2: amplitude of the steady-state vibration of the hydraulic system caused by $h_z(t) = h_{z0} \cos(\omega t)$

As derived in Lecture 7, slide 11, the amplitude of the steady-state response of the corresponding mechanical system is given as

$$U_{\text{steady}} = \frac{u_{\text{static}}}{\sqrt{\left(1 - \omega^2 / \omega_0^2\right)^2 + 4\zeta^2 \, \omega^2 / \omega_0^2}}, \qquad u_{\text{static}} = \frac{F_0}{k} = \frac{f_0}{\omega_0^2}$$

Therefore, using the relevant expressions for the hydraulic system (slide 13 of the current lecture), we obtain

$$H_{\text{steady}} = \frac{f_0}{\omega_0^2} \frac{1}{\sqrt{\left(1 - \omega^2 / \omega_0^2\right)^2 + 4\zeta^2 \,\omega^2 / \omega_0^2}}$$
$$\omega_0^2 = \frac{1}{BM}, \ \zeta = \frac{W}{2} \sqrt{\frac{B}{M}}, \ f_0 = \frac{h_{z0}}{BM}$$

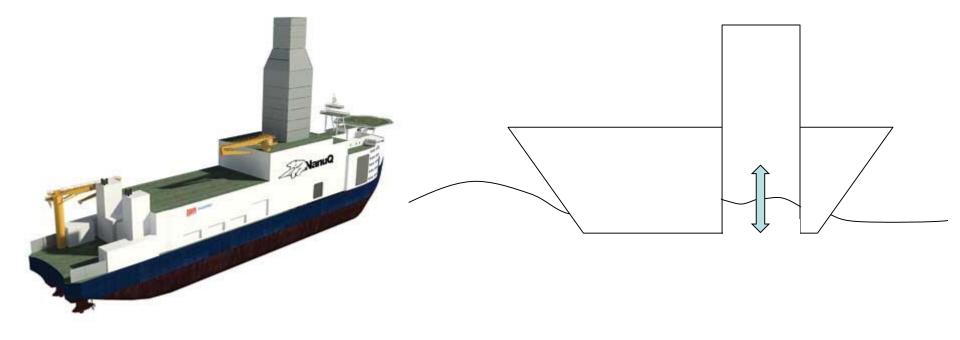
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Other examples of SDOF dynamical systems: pressure fluctuations in an enclosed moonpool of an arctic drillship



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Other examples of SDOF dynamical systems: the Helmholtz resonator (sound trapper)



$\omega_0 = c_{\text{sound}} \sqrt{\frac{A}{V_0 L}}$

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Other examples of SDOF dynamical systems: a loudspeaker



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Other examples of SDOF dynamical systems: voice production





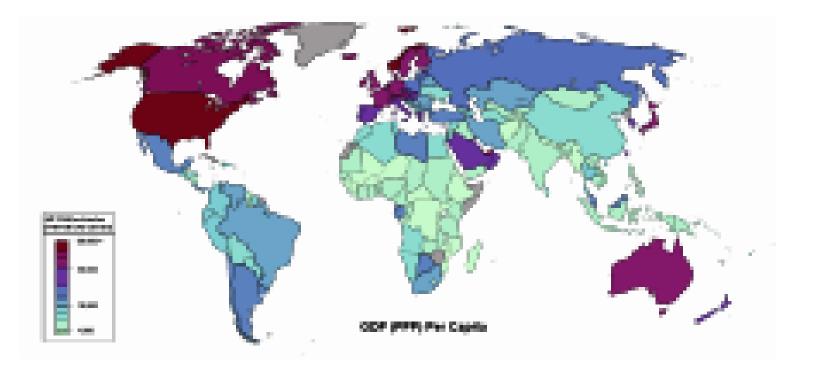
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Other examples of SDOF dynamical systems: business cycle



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Other examples of SDOF dynamical systems: predator-pray interaction

 $\dot{x} = bx - pxy \text{ (pray)}$

x(t) – Number of the prays (rabbits)

 $\dot{y} = -dy + rxy$ (predator) y(t) – Number of the predators (foxes)

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